

XProtoSphere: an eXtended multi-sized sphere packing algorithm driven by particle size distribution

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Abstract The sphere packing problem, which involves filling an arbitrarily shaped geometry with the maximum number of non-overlapping spheres, is a critical research challenge. ProtoSphere is a prototype-oriented algorithm designed for solving sphere packing problems. Due to its easily parallelizable design, it exhibits high versatility and has wide-ranging applications. However, the controllable regulation of particle size distribution (PSD) produced by ProtoSphere is often neglected, which limits its application on algorithm. This paper proposes a novel PSD-driven technique that extends the ProtoSphere algorithm to achieve multi-sized sphere packing with distribution-specific characteristics, as dictated by a pre-defined cumulative distribution function. The proposed approach improves the controllability and flexibility of the packing process, and enables users to generate packing configurations that meet their specific requirements. In addition, by combining the relaxation method with the ProtoSphere algorithm, we can further improve the packing density and ensure the average overlap below 1%. Our method generates multi-sized particles that can be used to simulate the behavior of various granular materials, including sand-like and clay-like soils.

Keywords ProtoSphere · Multi-Sized Sphere Packing · Particle Size Distribution · Discrete Element-Based Relaxation · Physically-Based Simulation

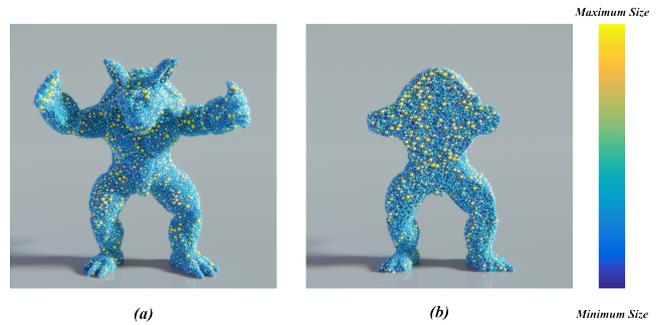


Fig. 1: Armadillo's XProtoSphere packing results (a), as well as corresponding cross-sectional views (b)

1 Introduction

The sphere packing algorithm has broad application to the process of filling particles densely within a given boundary without overlapping. It is commonly utilized to solve the problem of optimal sphere packing in three dimensions in mathematics [15], which is relevant to many fields such as coding theory and cryptography. In materials science, sphere packing algorithms aid in designing materials with desirable physical properties [29]. In wireless communication, they play a critical role in setting up the closest arrangement of antennas to optimize signal quality and coverage [9], and so forth [17].

In computer graphics community, sphere packing is frequently used for efficient spatial segmentation, collision detection [33], automatic rigging [3] and physically based simulation for granular materials [5,31]. In particular, granular material simulations require non-overlapping spheres for computational stability and a high packing density to simulate realistic sediment structures. Many of these applications utilize multi sized sphere packing algorithms rather than uniform sphere

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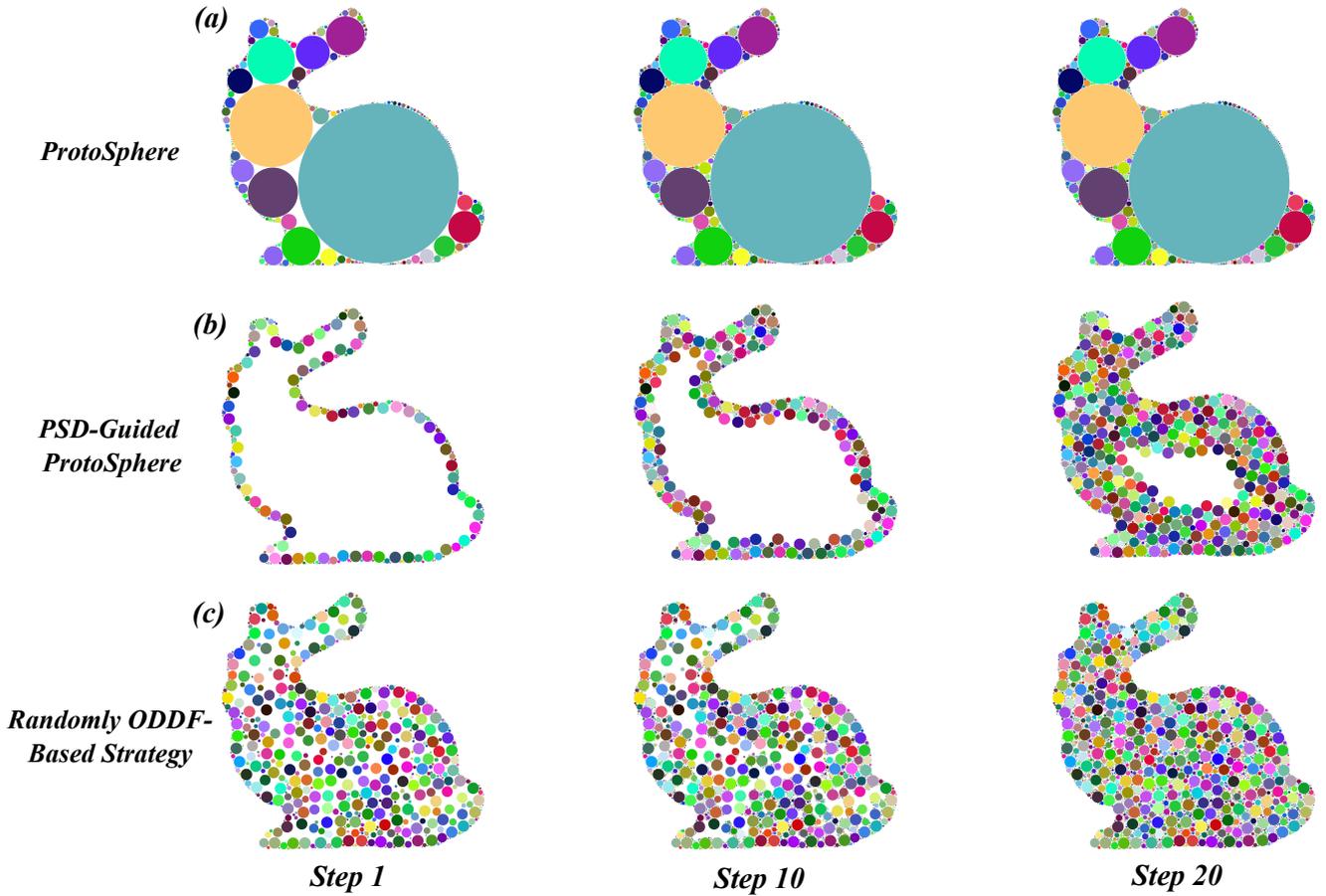


Fig. 2: Comparison of multiple particle insertion types based on the ProtoSphere algorithm, using a 2D Bunny polygon

50 packing algorithms, due to the flexibility and adaptabil-
 51 ity that multi-sized spheres offer in effectively model-
 52 ing complex physical systems with varying particle sizes
 53 and densities. In contrast, uniform sphere packing algo-
 54 rithms are often limited in their applicability, as they
 55 are primarily suitable for modeling homogeneous particle
 56 systems. Furthermore, uniform sphere packing can
 57 be easily achieved by fine-tuning certain sampling algo-
 58 rithms, such as the fast Poisson disk sampling method
 59 [8] or SPH-based blue noise sampling [21]. However,
 60 when these sampling algorithms are applied to multi-
 61 sized sphere packing problem, their efficacy in physi-
 62 cal simulations may be impeded by a higher overlap-
 63 ping rate and porosity within the sampling space [32].
 64 Therefore, these sampling methods may not be directly
 65 applicable for certain physical simulation applications,
 66 especially in the Discrete-Element Method (DEM) [11].

67 In comparison to the sampling algorithms and uni-
 68 form sphere packing algorithm mentioned before, Weller
 69 et al. [34] introduced a multi-sized packing algorithm
 70 called ProtoSphere. The ProtoSphere algorithm is in-

71 spired by the prototype-based approach in machine learn-
 72 ing, and it is capable of efficiently handling arbitrarily
 73 shaped objects. Furthermore, the algorithm is highly
 74 parallelizable, which makes it a promising option for a
 75 wide range of sphere packing problems. However, the in-
 76 ability of the ProtoSphere method to precisely control
 77 the particle size distribution of the generated spheres
 78 may limit its applicability, in particular granular simu-
 79 lation tasks, such as simulating soil structures contain-
 80 ing particles of varying sizes. This limitation may affect
 81 the ability to achieve more realistic simulation results.

82 To address these problems, this paper makes the
 83 following contributions:

- 84 – An extended algorithm that is based on the stan-
 85 dard ProtoSphere algorithm. The proposed algorithm
 86 enables users to predetermine a target particle size
 87 distribution using a cumulative distribution func-
 88 tion, thereby allowing for greater control over the
 89 particle size distribution of packed spheres.
- 90 – An randomly Offset Discrete Distance Field (ODDF)
 91 based strategy is proposed for achieving faster con-

vergence of particle size, as well as addressing the issue of boundary expansion towards the center that arises during particle generation using the extended ProtoSphere algorithm.

- A Discrete-Element based particle relaxation method is proposed to improve the packing density (see Figure 1). This method can be integrated with the extended ProtoSphere algorithm and is applicable in physically-based simulations. Compared to the SPH-based particle relaxation method, the approach offers greater stability and can be applied to a wider range of multi-sized particle distributions.

2 Related Work

The focus of this review section is to investigate the research and applications of the sphere packing problems in computer graphics and related fields. As is commonly understood, the proper placement of particles plays a vital role in the study of particle-based physically simulation animation. Particle-based methods typically use uniform particles, although in certain circumstances, such as those involving adaptive methods [12, 1, 35, 36] or DEM-related frameworks [30, 31], the application of multi-sized particles becomes imperative. With the advantage of being adaptable to arbitrary dimensions and objects, the ProtoSphere method was extended in our study to enable enhanced manipulation of particle size distribution. The resulting approach was explored for its suitability in physically-based simulation scenarios, specifically those involving DEM.

2.1 Sampling Based Approaches

In terms of the sampling and packing problems, although their objectives and application scenarios may differ, both algorithms are fundamentally concerned with the distribution of points in space. Thus, the points generated by the sampling method can be employed as centers for particles in a sphere packing algorithm by assigning each point a suitable radius.

Poisson Disk Sampling Blue noise sampling is a widely used technique in computer graphics due to its ability to produce a uniform distribution [38], making it applicable in a wide variety of applications [40, 22, 37, 20]. Poisson disk sampling, one of its patterns, is widely employed in rendering, geometry processing, and physically-based simulation due to its numerous applications and versatility. In particular, the faster versions of Poisson disk sampling, enhanced by Bridson [8], exhibit greater adaptability to arbitrary dimensions and

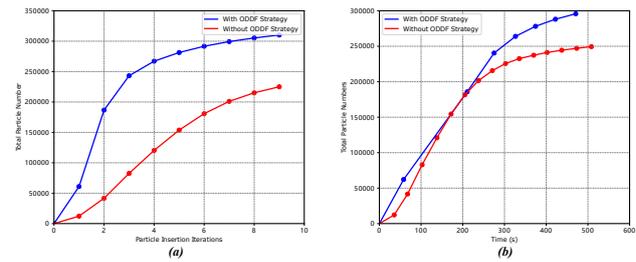


Fig. 3: Comparison of particle packing results obtained with and without the randomly offset discrete element field strategy demonstrates the efficiency of this technique in achieving higher packing densities

are comparatively easier to implement. Although this method has demonstrated commendable performance and broad applicability, regulating the packing density of the resulting particles can be challenging, especially in the case of multi-sized particle packing, where particles may overlap.

SPH Based Relaxation In general, the process of sampling particles for a given boundary involves dividing the plane into uniform grids in 2D (or voxelizing in 3D) and generating a particle within each uniform grid [26]. However, this method may produce an irregular distribution of particles near the boundary, which has the potential to obscure the original geometry's boundary information. To overcome this issue, Schechter et al. [25] employed a Poisson disk relaxation method to facilitate surface and volume sampling. Subsequently, Jiang et al. [21] attempted to use a cohesion-term integrated SPH method for blue noise sampling, which yielded promising outcomes for the relaxation of boundary particles. Moreover, they demonstrated that their method can be combined with adaptive methods [1] to facilitate multi-sized particle sampling. However, these approaches are all based on the SPH-based particle sampling method, and none of them are able to avoid the issue of large overlap between particles. Another problem remains that when particles are not uniformly distributed within the SPH kernel, these algorithms may become unstable.

2.2 Multi-Sized Sphere Packing

The algorithms that relate to multi-sized sphere packing can be categorized into three principal groups, namely geometry separation-based, mesh-based, and Apollonian-based methods:

Geometric Separation Based Approach Geometric separation-based algorithms focus on the task of

174 packing multi-sized particles by randomly inserting par-
 175 ticles within regions with low filling rate and itera-
 176 tively removing any overlapping particles. As an ex-
 177 ample, Lopes et al. [23,24] proposed a two-dimensional
 178 geometric separation method that enables the control of
 179 both porosity and particle size distribution through the
 180 use of a grid mapping approach. This method achieves
 181 high-efficiency particle insertion and removal, thereby
 182 facilitating the packing of multi-sized spheres. While ge-
 183 ometric separation-based algorithms can control poros-
 184 ity and particle size distribution, their strong stochas-
 185 ticity and the possibility of repeated insertion and re-
 186 moval of particles can make it difficult to ensure their
 187 high performance and low error rate when extended to
 188 3D space.

189 **Mesh Based Sphere Packing** The objective of these
 190 investigations is primarily to develop particle-based, non-
 191 overlapping geometries for use in DEM methods. The
 192 mesh-based methods utilized involve triangulating (2D)
 193 [4] or tetrahedral partitioning (3D) [14] in a given do-
 194 main, where particles can be positioned at vertices or
 195 within the unit geometry [10,18,19]. Recently, a re-
 196 fined approach for multi-sized particle packing has been
 197 proposed by Zhang et al. [39], which enables efficient
 198 and precise packing of particles for arbitrary 2D ge-
 199 ometries. This is achieved by improving upon Cui et
 200 al.’s algorithm [10] and utilizing a strategy that in-
 201 volves placing particles at each vertex of the triang-
 202 ular surface. Wang et al. [32] presented a novel method
 203 for optimizing porosity to enhance packing density us-
 204 ing the Power diagram [2]. This approach allows for
 205 the predefinition of particle size distribution, but the
 206 outcomes attained by this method suffer from an error
 207 rate of 10-20%. Both mesh-based studies face challenges
 208 when attempting to extend their algorithms to three-
 209 dimensional spaces, due to issues with performance and
 210 algorithm instability.

211 **Apollonian Based Method** The Apollonian packing
 212 algorithm [7] necessitates the initial placement of three
 213 mutually tangent discs, with each disc touching the
 214 other two. Subsequently, the algorithm iteratively in-
 215 serts additional discs into the largest available circular
 216 cavity within the remaining gap, utilizing this process
 217 to generate fractals of arbitrary dimensionality. The
 218 ProtoSphere method [34] is inspired by the prototype-
 219 guided approach in machine learning and employs an
 220 optimization process that utilizes multi-sized particles
 221 to fill geometries of arbitrary dimensions. The method
 222 yields results comparable to those produced by Apollo-
 223 nian sphere packing, while also being capable of accom-
 224 modating arbitrary geometries that are challenging to

Algorithm 1 Parallel ProtoSphere Algorithm

Input: surface Ω of object O , required particle number N
Output: a group of particles with radius information
 1: $\mathcal{D}_\Omega \leftarrow$ initialize the discrete distance field
 2: **repeat**
 3: $S : \{\mathbf{p}_1, \mathbf{p}_2, \dots, \mathbf{p}_n\} \leftarrow$ place prototype \mathbf{p}_i randomly
 inside grid c_i
 4: **for each** \mathbf{p}_i in S **do**
 5: **repeat**
 6: $\mathbf{q}_c = \arg \min \{\|\mathbf{p}_i - \mathbf{q}\| : \mathbf{q} \in \Omega\}$
 7: $\mathbf{p}_i \leftarrow \mathbf{p}_i + \varepsilon(t) \cdot (\mathbf{p}_i - \mathbf{q}_c)$
 8: $r_i = \|\mathbf{p}_i - \mathbf{q}_c\|$
 9: **until** \mathbf{p}_i has converged
 10: **end for**
 11: sort P by max radius r_i
 12: find $\mathbf{p}_k \in P$ that are not overlapped by any \mathbf{p}_i
 13: insert particles at positions \mathbf{p}_k with radii r_k
 14: update discrete distance field \mathcal{D}_Ω by $\Omega \leftarrow \Omega \cup \Omega_{\mathbf{p}_k}$
 15: **until** number of inserted particles $> N$

achieve using the latter method. Subsequently, Teuber 225
 et al. [28] proposed a GPU-based adaptive grid method 226
 that significantly enhances the performance of Proto- 227
 Sphere. Recently, Bonneau et al. [6] sought to incorpo- 228
 rate the multi-sized particles generated by ProtoSphere 229
 into the DEM, and achieved this by constraining the 230
 random point locations to control the size of the packed 231
 particles within a range pre-defined by the user. How- 232
 ever, currently available Apollonian-based methods fall 233
 short in achieving optimal particle size distribution. To 234
 address this limitation, this paper proposes an exten- 235
 sion of the ProtoSphere method that enables precise 236
 management of the particle size distribution. Addition- 237
 ally, the performance of this extended approach is eval- 238
 uated in the context of a physically-based particle sim- 239
 ulation. 240

3 ProtoSphere 241

Weller et al. [34] introduced the ProtoSphere algorithm,
 which is centered around the determination of the sphere
 radius by measuring the shortest distance between a
 point and the surface. To be specific, let Ω represent
 the surface of an arbitrary object O . The point \mathbf{q}_c on
 surface Ω that is closest to point \mathbf{p} can be defined as
 follows:

$$\mathbf{q}_c = \arg \min \{\|\mathbf{p} - \mathbf{q}\| : \mathbf{q} \in \Omega\} \quad (1)$$

Here, point \mathbf{p} can represent any position located 242
 within the interior of object O . The generated parti- 243
 cle is centered at point \mathbf{p} and has a radius $\|\mathbf{p} - \mathbf{q}_c\|$. 244

To approximate Apollonian-like sphere packing, the
 ProtoSphere algorithm employs a prototype-guided strat-
 egy that considers point \mathbf{p} as a prototype and seeks to

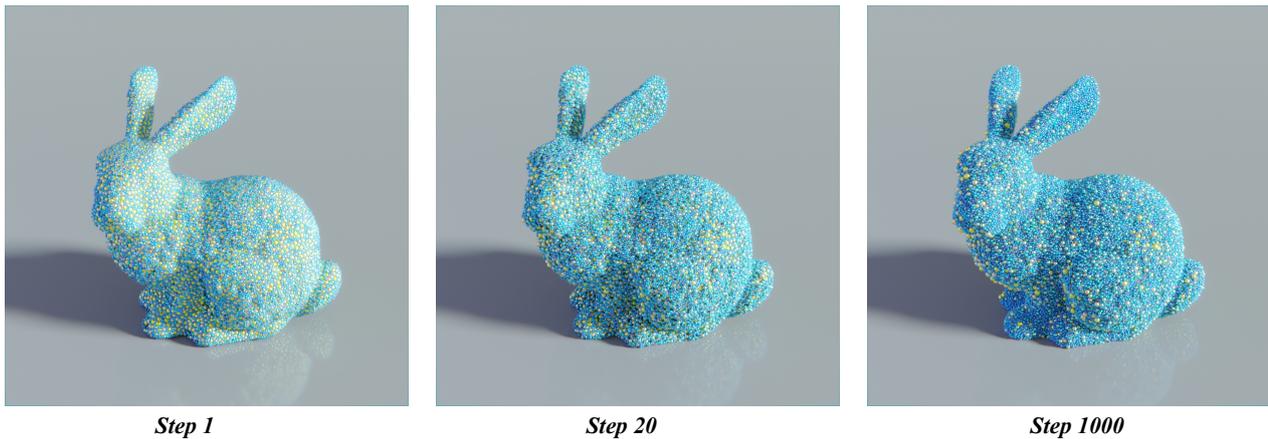


Fig. 4: Visualization of the particle movement process during coupling with the Discrete Element Relaxation

maximize its distance from the surface Ω (thus obtaining a particle with the largest possible radius within object O). The described process can be achieved using the following equation:

$$\mathbf{p} \leftarrow \mathbf{p} + \varepsilon(t) \cdot (\mathbf{p} - \mathbf{q}_c) \quad (2)$$

where the cooling function $\varepsilon(t) \in [0, 1]$ is employed to facilitate large movements during initial iterations and finer adjustments in subsequent steps. By employing this approach, the standard version of the ProtoSphere algorithm can be implemented through the following three steps: 1. Stochastically generate a point \mathbf{p} that lies within the interior of object O . 2. Update the position of point \mathbf{p} using Equation(2) iteratively until the convergence criterion is met. 3. Insert a particle with radius $\|\mathbf{p} - \mathbf{q}_c\|$ at point \mathbf{p} and return to the first step. It should be noted that the addition of each new particle necessitates the availability of surface information Ω_p to facilitate the updating of surface $\Omega = \Omega \cup \Omega_p$.

Importantly, the requirement for surface information with each new particle may lead to computational performance issues. Moreover, the standard ProtoSphere algorithm is confined to local optimization for each prototype and does not achieve global optimization. To overcome these limitations, Weller et al. [34] employs a gridding strategy that partitions the interior space of object O . With this approach, a prototype can be positioned within each grid, enabling global optimization by allowing them to move independently. In addition, they enhance the computational efficiency of determining the nearest boundary point by pre-computing the discrete distance field. The parallel version of the ProtoSphere algorithm is presented in Algorithm 1 with a detailed description. Figure 2(a) showcases a two-dimensional result obtained after computation with the parallelized algorithm.

4 Extended ProtoSphere (XProtoSphere)

ProtoSphere presents a superior option in sphere packing algorithms, as it possesses the capability to address arbitrarily shaped and multi-dimensional particle packing challenges. Additionally, its algorithmic implementation is straightforward and lends itself to parallelization, further increasing its competitiveness. Nevertheless, the fractal characteristics of the results generated by ProtoSphere pose challenges to its direct utilization in DEM-related granular simulations. Bonneau et al. [6] made an attempt to constrain the size of the particles generated by ProtoSphere by setting a range limit. However, they acknowledged explicitly that their method does not provide a means to regulate the particle size distribution of the outcomes. Therefore, our extended approach endeavors to enable control over the particle size distribution and address associated challenges.

4.1 PSD-Guided ProtoSphere

Modeling granular materials requires the determination of the relative proportion of particles of different sizes present in the material, which can be characterized by a piecewise constant distribution function $f(r)$ according to the following expression:

$$f(r) = \begin{cases} P_0 & \text{if } r_0 \leq r < r_1 \\ P_1 & \text{if } r_1 \leq r < r_2 \\ \vdots & \vdots \\ P_{n-1} & \text{if } r_{n-1} \leq r < r_n \end{cases} \quad (3)$$

where the probability density function (PDF) of each particle size interval $[r_0, r_1), [r_1, r_2), \dots, [r_{n-1}, r_n)$ is a constant value P_0, P_1, \dots, P_{n-1} .

297 To achieve a controlled particle size distribution,
 298 it is essential to ensure that the radii of the parti-
 299 cles generated by ProtoSphere align as closely as possi-
 300 ble with the PDF $f(r)$. Assuming that a set of target
 301 radii that conform to the PDF $f(r)$ can be obtained
 302 beforehand, the objective is then turn to ensure that
 303 the radii of all the particles generated by ProtoSphere
 304 converge to their respective target radii. Regarding the
 305 pre-calculation of the target radii, the cumulative dis-
 306 tribution function (CDF) $F(r)$ of the PDF $f(r)$ can be
 307 computed by $F(r) = \int_{-\infty}^r f(t)dt$ [13].

The fundamental concept behind our approach is to assign a target radius r'_i to each prototype \mathbf{p} during the gridding process (in Algorithm 1 line 3), and then to iteratively adjust the radius of each prototype r_i until it converges to its target radius r'_i . We can modify Equation(2) as follows:

$$\mathbf{p} \leftarrow \mathbf{p} + \varepsilon(t) \cdot (r' - r) \frac{\mathbf{p} - \mathbf{q}_c}{\|\mathbf{p} - \mathbf{q}_c\|} \quad (4)$$

This modified formula indicates that if the target radius r' exceeds the current radius $r = \|\mathbf{p} - \mathbf{q}_c\|$, point \mathbf{p} will move away from \mathbf{q}_c , otherwise it will move towards \mathbf{q}_c if the current radius is greater than the target radius. For the cooling function $\varepsilon(t)$, we found that a time-based decay function performed better in our experiments. The function is shown as follows:

$$\varepsilon(t) = \frac{\varepsilon(t-1)}{1+kt} \quad (5)$$

308 where k is a parameter that regulates the decay rate,
 309 and we set its value to 0.01 for all experiments.

310 4.2 Randomly ODDF-Based Strategy

311 In Section 4.1, we introduce a modification to the pro-
 312 totype's motion, whereby its radius continuously con-
 313 verges towards the predetermined radius. Nevertheless,
 314 in contrast to the standard ProtoSphere algorithm, our
 315 approach yields relatively diminutive particle sizes in
 316 the initial stages of particle insertion. This results in
 317 the initial particles being inevitably placed at the bound-
 318 ary of the object, and subsequent inserted particles
 319 gradually expand from the boundary towards the center,
 320 which is an undesirable outcome. Specifically, this
 321 method of particle insertion gives rise to two main is-
 322 sues. Firstly, inserting particles in this way requires
 323 more steps, as each newly inserted particle can be seen
 324 as a new boundary of the object. Secondly, it results
 325 in a more homogeneous distribution of the inserted par-
 326 ticles, as the particle size distribution of the newly in-
 327 serted particles should be similar for each layer. Figure
 328 2(b) illustrates the process of particle insertion using

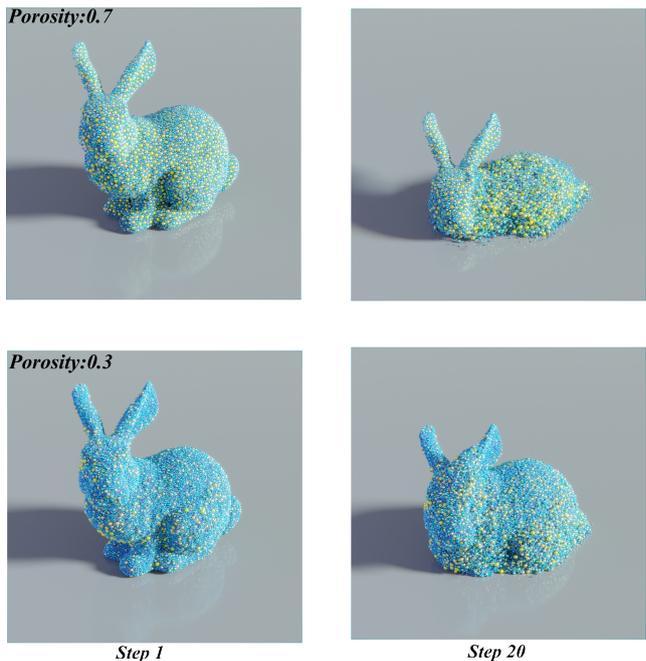


Fig. 5: Comparison experiment on the effects of multi-sized particle packing on capillary forces at two porosities

the 2D PSD-Guided ProtoSphere method, where it is evident that the process starts from the boundary and gradually progresses towards the center.

This issue arises due to the standard ProtoSphere algorithm's determination of the radius for each particle to be inserted, which is based on the distance between the current particle \mathbf{p} and the nearest point \mathbf{q}_c on the boundary. More specifically, the standard ProtoSphere algorithm neglects the consideration of particle radius sizes, leading to the insertion of several large-radius particles during the initial stages of particle insertion, forming the internal skeleton of the geometry. While our modified ProtoSphere method (Equation(4)) is designed to regulate the size of the particle radius, and in the majority of cases, it is unnecessary to generate particles with such a large radius. Consequently, the process of gradually inserting particles layer by layer from the boundary towards the center, as illustrated in Fig-

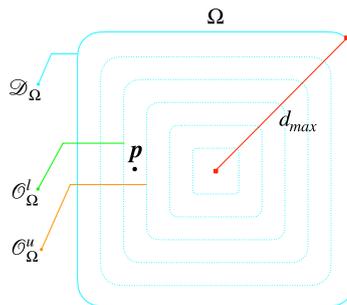


Fig. 6: Illustration of the randomly offset discrete distance field-based strategy for particle packing

ure 2(b), is observed. To address the issue of undesired particle insertion, our proposed solution involves assigning a unique discrete distance field to each particle. This allows for the recalculation of both the displacement direction $\mathbf{p} - \mathbf{q}_c$ and the particle radius r . We achieve this by introducing a random variable $\rho \in [0.04, 0.2]$ to each prototype. Then we divide the discrete distance field \mathcal{D}_Ω into multiple subfields based on the maximum distance d_{\max} in the field. As a result, we obtain multiple distance fields with varying offset levels, as shown by the dashed lines in Figure 6. Based on the position of current prototype \mathbf{p} in the distance field, we locate the two offset distance fields, \mathcal{O}_Ω^u and \mathcal{O}_Ω^l that are closest to it. At last, the updated distance field \mathcal{D}'_Ω can be computed by utilizing the pre-computed distance field \mathcal{D}_Ω , the offset distance fields \mathcal{O}_Ω^u and \mathcal{O}_Ω^l , as follows:

$$\begin{aligned} \mathcal{D}'_\Omega(\mathbf{p}) &= \min(\mathcal{D}_\Omega(\mathbf{p}), \min(\mathcal{O}_\Omega^u(\mathbf{p}), \mathcal{O}_\Omega^l(\mathbf{p}))) \\ \mathcal{O}_\Omega^u(\mathbf{p}) &= \left\lceil \frac{\mathcal{D}_\Omega(\mathbf{p})}{\rho d_{\max}} \right\rceil \rho d_{\max} - \mathcal{D}_\Omega(\mathbf{p}) \\ \mathcal{O}_\Omega^l(\mathbf{p}) &= \mathcal{D}_\Omega(\mathbf{p}) - \left\lfloor \frac{\mathcal{D}_\Omega(\mathbf{p})}{\rho d_{\max}} \right\rfloor \rho d_{\max} \end{aligned} \quad (6)$$

As illustrated in Figure 2(c), our proposed strategy can insert particles of the target size at unpredictable locations in the space during the initial insertion phase. This is in contrast to Figure 2(b), where particle insertion is constrained to the boundaries only. Moreover, we evaluated the performance of the XProtoSphere algorithm with and without the randomly ODDF strategy, and obtained promising results. As shown in Figure 3(a), the algorithm with the ODDF strategy is capable of inserting more particles at each iteration, resulting in a higher total number of inserted particles compared to the non-ODDF strategy. To further evaluate the algorithm's efficiency, we measured the computation time for the ODDF and non-ODDF strategies, as depicted in Figure 3(b). The results demonstrate that the ODDF strategy consistently outperforms the non-ODDF strategy in generating the same number of particles. This suggests that although the ODDF strategy requires more computational resources, it is more efficient for particle insertion.

4.3 Coupling with Discrete Element Relaxation

In the standard ProtoSphere algorithm, the porous regions within the space can be continuously explored to identify and fill them with particles of appropriate sizes. However, when there is a limit on the particle size that we need to insert, it becomes challenging to maximize the packing density of the entire region and minimize

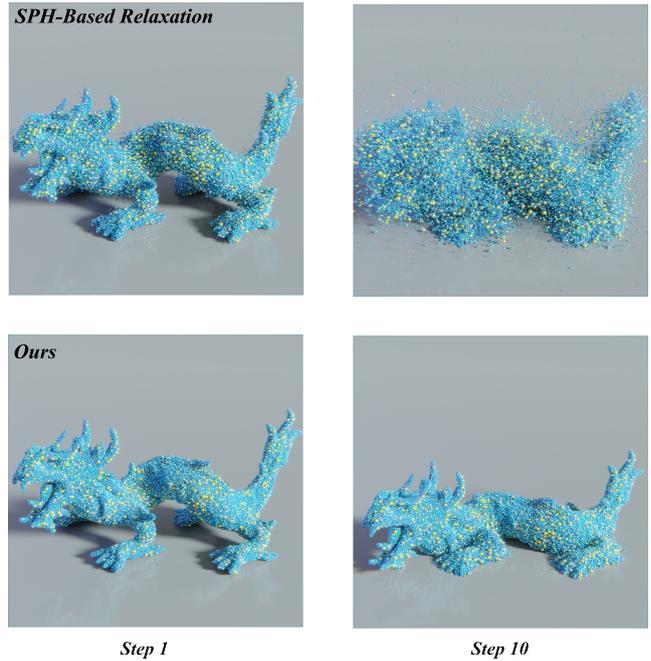


Fig. 7: Comparison between multi-sized particles packing using the SPH-based relaxation and our method when applied to DEM

its porosity. It is essential to acknowledge that the inability to further reduce the spatial porosity is mainly attributed to the suboptimal distribution of the already inserted particles, which limits our capability of further particle insertion. If we induce particle movements, then more space can be created during the motion, thereby allowing for the insertion of additional particles. The ProtoSphere algorithm detects particle overlaps during the insertion of new particles and subsequently removes them. However, our proposed approach involves relaxing the strict constraints on particle overlap during the

Algorithm 2 XProtoSphere Coupling with Discrete Element Relaxation

Input: surface Ω of object O , a probability density function $f(r)$, maximum relaxation steps N

Output: a group of particles with radius information

- 1: insertion number $\ell \leftarrow 1$
 - 2: **repeat**
 - 3: Using XProtoSphere($\Omega, f(r)$) to insert particles once
 - 4: **for** relaxation steps $i = 0; i < N; i = i + 1$ **do**
 - 5: elastic force \mathbf{F}_k^n between the particles can be computed by employing Equation(7)
 - 6: compute the acceleration \mathbf{a} of particles and update it using the energy decay mechanism outlined in Equation(8)
 - 7: use Equation(9) corrects the position and velocity of particles near the boundary
 - 8: advance the particle
 - 9: **end for**
 - 10: **until** $\ell >$ maximum insertion number
-

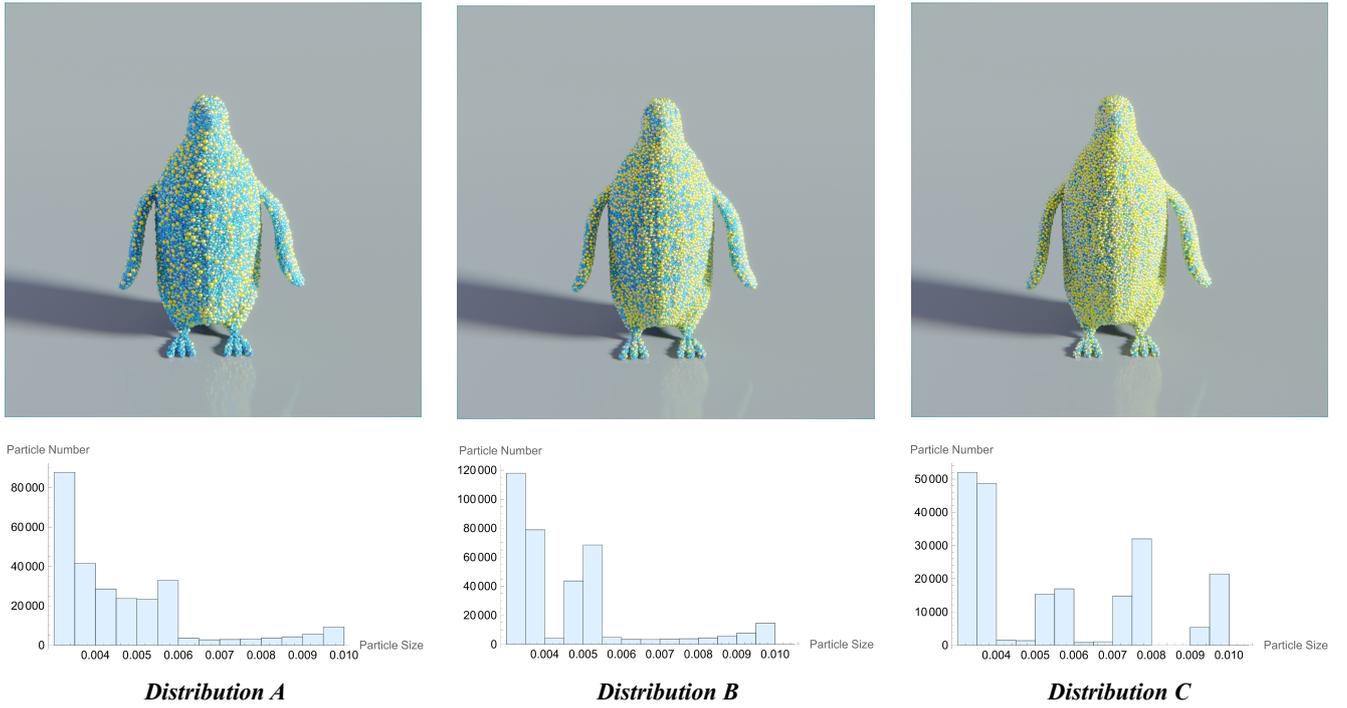


Fig. 8: Comparison experiment using XProtoSphere for multi-sized particle packing under different pre-defined particle size distributions

401 insertion phase. Subsequently, we utilize the DEM to
402 iteratively reduce the overlap between particles.

The DEM is a widely used Lagrangian-based simulation model for granular media, commonly employed in the field of soil mechanics [11]. In the DEM system, each particle is treated as a rigid body with a defined position and radius. When the distance between two particles, denoted as $d = \|\mathbf{x}_{k_1} - \mathbf{x}_{k_2}\|$, is less than the sum of their radii (i.e., $r_{k_1} + r_{k_2}$), they are considered to be in contact or collide with each other. At this point, normal and tangential forces, denoted as \mathbf{F}_k^n and \mathbf{F}_k^t , respectively, are exerted on the particles to simulate their interactions [16]. Our proposed approach focuses on utilizing only the elastic force \mathbf{F}_k^n in the normal direction to separate overlapping particles, define as:

$$\begin{aligned} \mathbf{F}_k^n &= K_N(r_{k_1} + r_{k_2} - d)\hat{\mathbf{x}}_{k_1k_2} \\ K_N &= \frac{2Y_1R_1Y_2R_2}{Y_1R_1 + Y_2R_2} \\ \hat{\mathbf{x}}_{k_1k_2} &= \frac{\mathbf{x}_{k_1} - \mathbf{x}_{k_2}}{d} \end{aligned} \quad (7)$$

403 where K_N represents the normal stiffness coefficient,
404 which is related to Young's modulus Y and particle
405 radius R .

To prevent continuous fluctuations caused by particles getting trapped in narrow gaps, we employ an energy decay mechanism [27] as follow:

$$\mathbf{a} \leftarrow \mathbf{a} \left(1 - \lambda \text{sgn} \left(\mathbf{a} \left(\mathbf{v} + \frac{\mathbf{a}\Delta t}{2} \right) \right) \right) \quad (8)$$

406 where the parameter λ represents the damping coefficient,
407 \mathbf{a} denotes acceleration, \mathbf{v} represents velocity, and
408 Δt is the time step. Additionally, the sgn function is
409 utilized with possible values of -1, 0, or 1.

It is important to note that when a particle begins to move due to elastic forces, it must be constrained to remain within the object. In particular, if the particle makes contact with the boundary, a repulsive force must be applied to prevent it from crossing the boundary. In our approach, if a particle \mathbf{p} exceeds the boundary Ω , we do not apply a repulsive force. Rather, we modify its position and adjust its velocity to ensure that it remains inside the object O . Specifically, we use position modification to move the particle back inside the boundary, and velocity adjustment to match the post-repulsive velocity, as shown below:

$$\begin{aligned} \mathbf{x} &\leftarrow \mathbf{x} - \mathcal{D}_\Omega(\mathbf{x})\nabla\hat{\mathcal{D}}_\Omega(\mathbf{x}) \\ \mathbf{v} &\leftarrow e(\mathbf{v} - 2(\mathbf{v} \cdot \hat{\mathcal{D}}_\Omega(\mathbf{x}))\hat{\mathcal{D}}_\Omega(\mathbf{x})) \end{aligned} \quad (9)$$

410 The equation involves the repulsive coefficient repre-
411 sented by the parameter e , and the normalized gra-
412 dient of the discrete distance field \mathcal{D}_Ω at position \mathbf{x} ,
413 denoted by $\hat{\mathcal{D}}_\Omega(\mathbf{x})$. Note that the formula presented
414 above is executed only when the distance $\mathcal{D}_\Omega(\mathbf{x})$ is less
415 than or equal to 0. The complete particle motion was
416 recorded during the application of the Discrete Element
417 Relaxation, as shown in Figure 4. The bunny-shaped
418 multi-sized spheres in Step 1 were generated using the

XProtoSphere, with a minor modification to the insertion process. Specifically, the insertion condition in line 18 of Algorithm 3 was altered to permit the insertion of spheres with an overlapping rate of ϵ and ξ or less, rather than those without any overlapping spheres. This is due to the fact that in the initial stages of the process, if the particles do not overlap, the Discrete Element Relaxation is unable to generate the necessary forces to move the particle swarm and adjust particle positions to optimize space filling. It is important to note that when coupled with the discrete element relaxation method, a small overlap between particles may still occur. Detailed experimental results on the spatial porosity and overlapping rate obtained by coupling XProtoSphere with discrete element relaxation are presented in Table 1.

4.4 Implementation Details

Algorithm 3 outlines the specific implementation details for XProtoSphere, highlighting the differences from the standard ProtoSphere algorithm in red. Compared to ProtoSphere, XProtoSphere necessitates a pre-defined probability density function $f(r)$ to set the target particle size distribution as input data. Regarding the randomly ODDF-strategy discussed in Section 4.2, it is worth noting that it is only necessary to apply this strategy during the initial particle insertion process (line 7). This is due to the fact that particles can already be inserted at non-boundary locations of the object O in the first operation, and the resulting inserted particles will subsequently be utilized as new boundary information in subsequent computations. During the particle insertion process, we apply a stricter constraint as indicated in line 18 of our algorithm.

In order to minimize the discrepancy between the target particle size and the actual size following insertion, we also require that $|\frac{r_k}{r'_k} - 1|$ not exceed the threshold $\epsilon = 0.01$. Typically, an acceptable particle overlapping rate (ξ which is described in Section 4.3) of 0.4 to 0.6 is employed, as this allows for minimal particle overlapping following implementation of the Discrete Element Relaxation. Algorithm 2 details the process of coupling XProtoSphere with Discrete Element Relaxation to further enhance the packing density. Upon the insertion of each particle by means of XProtoSphere, the iteration process is performed $N = 1000$ times using Discrete Element Relaxation.

5 Experimental Results

This section presents the results obtained by our algorithm for sphere packing of various geometries, along

Algorithm 3 XProtoSphere Algorithm

Input: surface Ω of object O , a probability density function $f(r)$

Output: a group of particles with radius information

- 1: insertion number $\ell \leftarrow 1$
- 2: $\mathcal{D}_\Omega \leftarrow$ initialize the discrete distance field
- 3: **repeat**
- 4: $S : \{\mathbf{p}_1, \mathbf{p}_2, \dots, \mathbf{p}_n\} \leftarrow$ place prototype \mathbf{p}_i randomly inside grid c_i
- 5: **for each** prototype \mathbf{p}_i , a target radius r'_i is generated independently based on $f(r)$
- 6: **if** $\ell == 1$ **then**
- 7: use Equation(6) to compute \mathcal{D}'_Ω as discrete distance field
- 8: **else**
- 8: update discrete distance field \mathcal{D}_Ω by $\Omega \leftarrow \Omega \cup \Omega_{\mathbf{p}_k}$
- 9: **end if**
- 10: **for each** \mathbf{p}_i in S **do**
- 11: **repeat**
- 12: $\mathbf{q}_c = \arg \min \{\|\mathbf{p}_i - \mathbf{q}\| : \mathbf{q} \in \Omega\}$
- 13: $\mathbf{p}_i \leftarrow \mathbf{p}_i + \epsilon(t) \cdot (r'_i - r_i) \frac{\mathbf{p}_i - \mathbf{q}_c}{\|\mathbf{p}_i - \mathbf{q}_c\|}$
- 14: $r_i = \|\mathbf{p}_i - \mathbf{q}_c\|$
- 15: **until** \mathbf{p}_i has converged
- 16: **end for**
- 17: sort P by max radius r_i
- 18: **find** $\mathbf{p}_k \in P$ that are not overlapped by any \mathbf{p}_i and $|\frac{r_k}{r'_k} - 1| < \epsilon$
- 19: insert particles at positions \mathbf{p}_k with radii r_k
- 20: $\ell = \ell + 1$
- 21: **until** $\ell >$ maximum insertion number

with its application in physically-based simulation algorithms. We also evaluated the performance of the XProtoSphere when applied to different geometries, as shown in Table 1. A, B, C meaning the particle size distributions shown in Figure 8.

The computational processes defined by lines 6 to 17 in Algorithm 3 and lines 4 to 8 in Algorithm 2 are executed on an NVIDIA GeForce RTX3090 24GB GPU by using CUDA programming framework. The pre-calculation of the discrete distance field and the computation of the particle insertion step are executed on a AMD Ryzen 9 3900X CPU.

Application to Capillary Model To accurately simulate sandy materials with clay-like behavior, achieving a higher packing density of multi-sized particles is crucial to enhance the realism of the simulation and distinguish the loose nature of the sand. In this study, we employed multi-sized particles with high porosity (low packing density) and low porosity (high packing density) to model sand with clay properties, as depicted in Figure 5. By applying the capillary force model proposed by Wang et al. [31], Our findings suggest that low-porosity particles exhibit a greater propensity for clay-like behavior compared to high-porosity particles. This finding shows that the proposed method is useful for reproducing sediment structures.

Table 1: Comparing XProtoSphere performance across different geometries

Model	Number			Time(s)			Overlap($\times 10^{-6}$)			Porosity(%)			BHD($\times 10^{-3}$)			JSD($\times 10^{-3}$)			KLD($\times 10^{-2}$)		
	A	B	C	A	B	C	A	B	C	A	B	C	A	B	C	A	B	C	A	B	C
Armadillo	243k	323k	174k	328	431	231	7.87	15.1	13.4	30	33.2	37.8	2.4	4.3	24	2.3	4.28	23.6	0.8	1.71	9.6
Beast	292k	401k	217k	410.1	592.56	315.53	8.63	23.39	12.18	29	32.8	35.7	3.8	5.5	24.13	3.7	5.5	23.7	1.27	2.2	9.6
Bunny	282k	387k	206k	384.53	571.01	278.3	8.519	10.7	12.4	31	35.2	38	3.5	6.12	27.9	3.43	6.09	27.44	1.17	2.4	11.1
Cheburashka	229k	324k	184k	256.249	367.13	203.9	8.1	11.6	11.3	28.7	34.3	34	6.34	10	26.6	6	9.9	26.1	2.06	4.04	10.6
Cow	303k	418k	226k	389.06	577.68	285.25	8.707	15.5	11.8	29.9	34.4	36.1	4.13	6.7	26.2	3.99	6.7	25.8	1.36	2.7	10.4
Dolphin	293k	415k	230k	363.56	555.11	280.4	7.888	32.1	10.5	28.8	33.8	34.7	5.66	7.5	26.3	5.44	7.4	25.9	1.84	3	10.5
Dragon	202k	274k	156k	221.17	286.7	180.47	11.44	15.5	33.1	28.9	35	33.5	5.9	15.6	27.1	5.7	15.4	26.6	1.952	6.3	10.8
Homer	221k	309k	174k	246.33	351.2	189	7.267	11.5	10.8	28.2	32.9	34.1	5.49	7.1	24.5	5.3	7	24.1	1.8	2.8	9.7
Horse	270k	371k	206k	353.29	508	273.3	8.508	11.9	12.7	30.5	34.6	36.2	4.4	7.1	26.1	4.3	7.1	25.7	1.46	2.8	10.4
Lucy	224k	306k	166k	400.41	349.95	200	10.25	11.8	13.5	29.3	34.5	35.2	4.67	10.2	26.3	4.5	10.2	25.8	1.53	4.1	10.4
Nefertiti	326k	443k	239k	502.9	740	373.75	16.2	92.9	12.6	31.3	33.7	38	2.94	4.1	24.9	2.8	4.1	24.5	0.99	1.6	10
Penguin	264k	364k	211k	299.99	417	232.5	8.027	12	10	28.9	34.5	33.4	6.2	11.6	28.6	5.9	11.5	28	2.02	4.7	11.5
Spot	292k	416k	232k	370.166	547.9	284.98	8.003	12.4	10	29.3	34.8	34.9	5.95	9	27.8	5.7	9.1	27.3	1.94	3.7	11.1
Suzanne	311k	448k	262k	404.163	606	337	7.97	13.7	10.9	33.6	39	37.8	6.8	11.3	33.5	6.5	11.3	32.8	2.23	4.6	13.4
Turtle	236k	347k	223k	267.365	386.6	260.46	60.77	12.2	152	25.57	33.6	29.6	13.6	18.4	30.8	12.9	18.1	30	4.39	7.3	12.9
Average	266k	370k	207k	346.485	485.856	261.722	12.543	20.152	22.478	29.531	34.42	35.266	5.452	8.968	26.982	5.23	8.911	26.489	1.787	3.596	10.8

494 **Comparison with SPH-Based Relaxation** To enhance packing density, a comparison experiment between the SPH-based relaxation [21] and our method was conducted using the DEM. Results from Figure 7 revealed that the SPH-based relaxation exhibited considerable instability during the early stages of simulation, which is likely due to high particle overlap resulting from the algorithm’s application to irregularly distributed multi-sized particles. In contrast, our proposed method demonstrates significantly more stable simulations, and increases packing density with minimal particle overlap, as shown in Table 1. Furthermore, we performed a quantitative analysis to assess the stability of our method by comparing the total energy, as shown in Figure 9. The blue line in the figure depicts the total energy variation of the DEM simulation with non-overlapping particles in the initial stage. Its variation trend is highly similar to that of our proposed method (green line) during the DEM simulation. However, the SPH Relaxation-based method produces abnormally high total energy values during the early stages of the simulation. The unstable simulation results, as shown in Figure 7, provide further evidence of this issue. As outlined in Section 4.3, our proposed method for further enhancing the particle packing density cannot completely eliminate particle overlap, leading to a slightly higher total energy in the DEM simulation compared to the case without particle overlap. Nevertheless, the average overlap rate between particles packed by our method is significantly below 1%, whereas that of the SPH-based method is approximately 30%. This discrepancy is also why our method can produce stable DEM simulations, while the SPH-based method cannot.

528 **Comparison Experiments on Various Predefined Particle Size Distributions** To illustrate the versatility of our method, we conducted experiments using three distinct representative particle size distributions, as shown in Figure 8. We present visualizations of the resulting multi-sized particle packing and corresponding particle distributions. For the quantitative analysis

535 results of these three sets of experiments, please refer to the data row labeled ‘Penguin’ in Table 1. 536

537 **Application to Discrete Element Method** Herein, we present the results of multi-sized particle packing experiments on 15 commonly used 3D models to demonstrate the applicability and versatility of our algorithm in handling objects of arbitrary geometry. The obtained packed particles are shown to be stable when applied to DEM-related algorithms for sand-like simulation (see Figure 10). In Table 1, we provide a comprehensive record of the results of our multi-sized particle packing experiments for the 15 commonly used 3D models, including the total number of particles, the computational time, the average overlapping rate, and the porosity. Moreover, we employed three metrics, Bhattacharyya Distance (BHD), Jensen-Shannon Divergence (JSD), and Kullback-Leibler Divergence (KLD), respectively, to quantitatively evaluate the difference between the particle size distribution obtained by our algorithm and the pre-defined particle size distribution. Specifically, these three indicators can serve as measures to assess the similarity between two distributions. A value closer to 0 indicates a higher level of similarity between the two distributions, and a value of 0 suggests that the distributions are identical. As shown by the results in Table 1 from the experiments conducted using three different predefined distributions, the indicators used to evaluate the similarity of the distributions range from approximately 0.1 to 0.001. These results highlight the effectiveness of our algorithm in generating particle distributions that closely resemble the pre-defined distributions. 566

567 6 Conclusion

568 This study presents the XProtoSphere method, a multi-sized particle packing algorithm that can conform to a pre-defined distribution by extending the ProtoSphere approach. Specifically, the radius of each prototype is gradually adjusted to match its target radius. Additionally, we propose a randomly ODDF-based strategy to 573

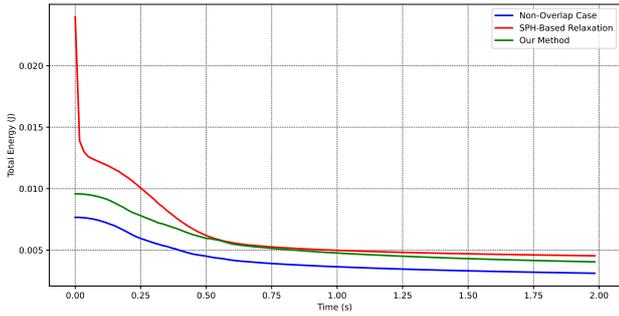


Fig. 9: Quantitative assessment of the stability of particle relaxation methods based on SPH and DEM through comparative total energy analyses

574 enhance the initial particle insertion phase and enable
 575 the packing of a larger number of particles. Finally, we
 576 integrate the XProtoSphere with the Discrete Element
 577 Relaxation method to further increase packing density
 578 and minimize particle overlap, which are essential for
 579 stable granular material simulations.

580 The major limitation of our algorithm that we need
 581 to remove overlapping particles and insert newly gen-
 582 erated particles without overlap when all prototypes
 583 converge. It presents two challenges: firstly, it is diffi-
 584 cult to parallelize this process on the GPU, as it may
 585 consume significant computational resources. Secondly,
 586 the removal of overlapping particles can result in a
 587 large discrepancy between the final packed particle dis-
 588 tribution and the pre-defined particle distribution. To
 589 address these issues, we would like to propose a fully
 590 GPU-based XProtoSphere algorithm in the further re-
 591 search, which can control particle size distribution more
 592 effectively by avoiding the need for particle deletion. In
 593 addition to computational performance-related issues,
 594 while our proposed method allows for the regulation of
 595 the particle radius distribution, it does not have the
 596 capability to regulate the spatial distribution of parti-
 597 cles with varying sizes. One potential solution involves
 598 utilizing the medial axis-based local geometric feature
 599 function introduced in Adams et al. [1] to regulate the
 600 spatial distribution of particles. However, this approach
 601 presents a challenge in simultaneously regulating the
 602 distribution of particle sizes. Therefore, in future work,
 603 we would like to propose a method that can effectively
 604 control both the particle size distribution and the spa-
 605 tial distribution of multi-sized particles. By doing so,
 606 we can extend the applicability of the XProtoSphere
 607 method to a wider range of scenarios.

608 Conflict of Interest Statement

609 The authors have no competing interests to declare that
 610 are relevant to the content of this article.

611 Data Availability Statement

612 Data will be made available on reasonable request.

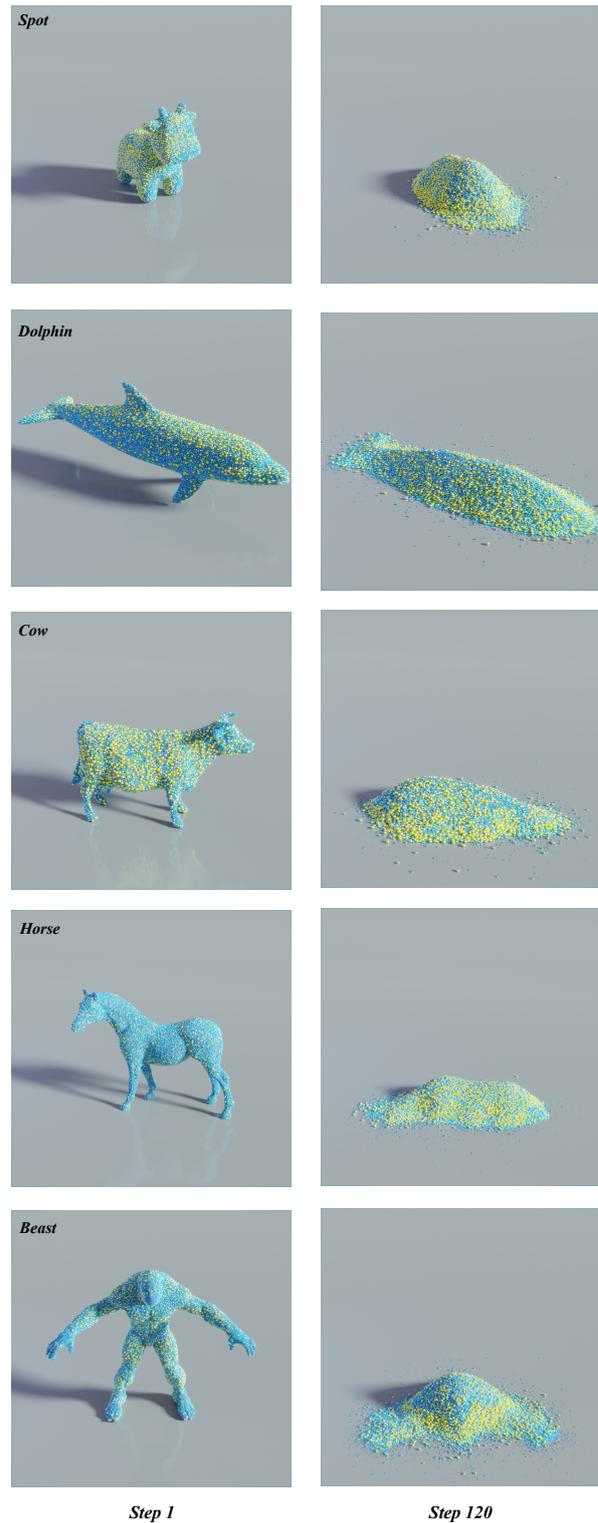


Fig. 10: Additional results generated by the XProtoSphere and their effects after importation into the sand simulator

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